

# Geometric Algebra: It's not Algebraic Geometry!

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# Geometric algebra...

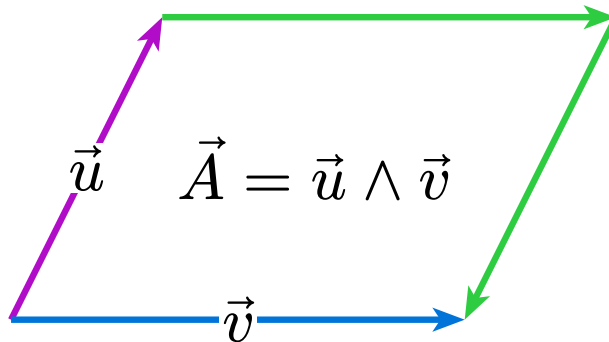
- Makes things intuitive
- Generalizes “weird” phenomena
  - Cross product
  - Quaternions for rotation
  - A lot of physics
- Provides a more powerful framework than linear algebra
  - You can add scalars and vectors (sort of)!

# Geometric Primitives

For now, consider 2D space. Let  $\vec{u} = (1, 2)^T$  and  $\vec{v} = (3, 0)^T$ :

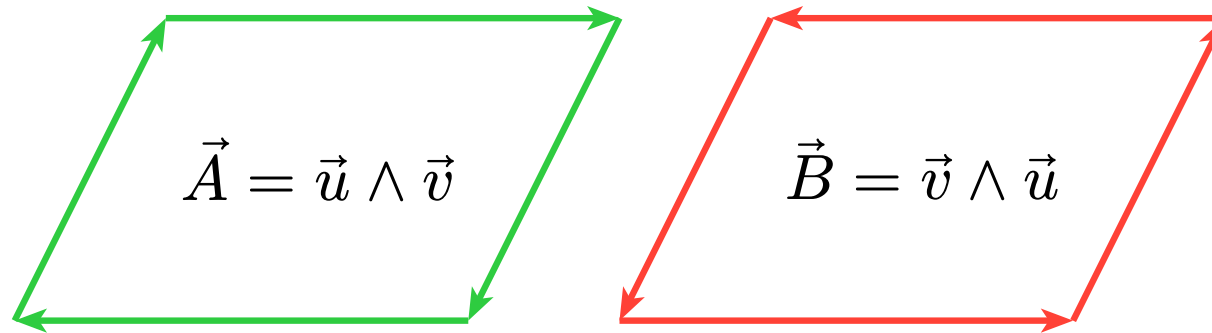


Then, we define their **wedge product**,  $\vec{A} := \vec{u} \wedge \vec{v}$ :

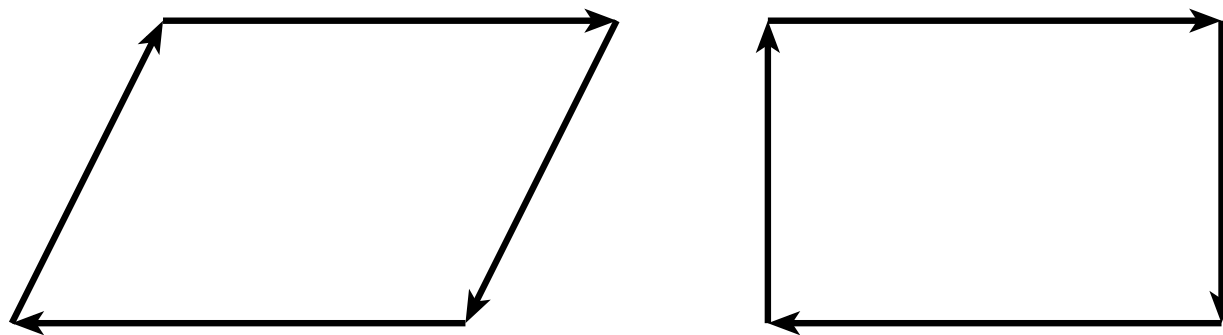


as the oriented 2D area formed by joining the two vectors. We call  $\vec{A}$  a **bivector**. Note that  $\vec{A}$  has an area and an orientation.

The wedge product is **anti-commutative**, i.e.  $\vec{u} \wedge \vec{v} = -\vec{v} \wedge \vec{u}$ . This can be visualized as the orientation of the two bivectors these wedge products form:



In fact, the wedge product is uniquely determined by the magnitude and orientation. The following two bivectors are equivalent:



In other words, shape does not matter.

## More Intuition for Wedge Product

We can infer two more properties about the wedge product:

1. The wedge product is **bilinear**.
2. For parallel  $\vec{u}, \vec{v}$ , we have that  $\vec{u} \wedge \vec{v} = 0$ .

Hence, let  $e_1, e_2$  be the standard basis vectors of  $\mathbb{R}^2$ . We can express the wedge product of  $\vec{u}, \vec{v} \in \mathbb{R}^2$  in terms of these basis vectors:

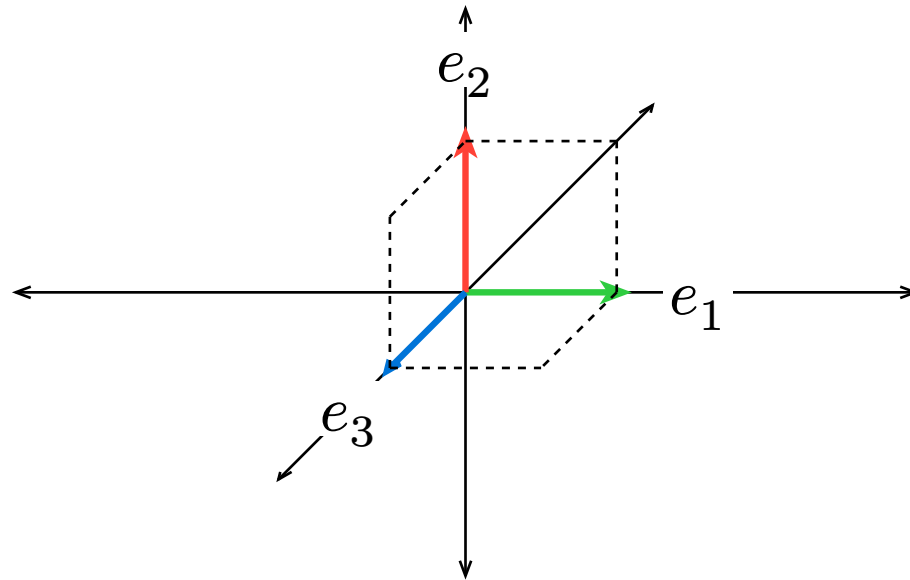
$$\begin{aligned}
 \vec{u} \wedge \vec{v} &= (\vec{u}_1 e_1 + \vec{u}_2 e_2) \wedge (\vec{v}_1 e_1 + \vec{v}_2 e_2) \\
 &= \vec{u}_1 \vec{v}_2 (e_1 \wedge e_2) + \vec{u}_2 \vec{v}_1 (e_2 \wedge e_1) \\
 &= \vec{u}_1 \vec{v}_2 (e_1 \wedge e_2) - \vec{u}_2 \vec{v}_1 (e_1 \wedge e_2) \\
 &= (\vec{u}_1 \vec{v}_2 - \vec{u}_2 \vec{v}_1) (e_1 \wedge e_2)
 \end{aligned}$$

Since  $e_1 \wedge e_2$  forms the unit square, we can see that the area of  $\vec{u} \wedge \vec{v}$  equals  $B_{12} := \vec{u}_1 \vec{v}_2 - \vec{u}_2 \vec{v}_1 = \sin(\theta) \|u\| \|v\|$ . Furthermore,  $\vec{u} \wedge \vec{v}$  form the unit square scaled by  $B_{12}$ .

**2D  $\rightarrow$  3D**

For now, bivectors are analogous to planes. In 3D space, bivectors have three components, instead of one. Let  $e_1, e_2, e_3$  be the standard basis of  $\mathbb{R}^3$ .

Vectors and bivectors in 3D:



$$\vec{v} = v_1 e_1 + v_2 e_2 + v_3 e_3$$

$$\vec{B} = B_{12}(e_1 \wedge e_2) + B_{13}(e_1 \wedge e_3) + B_{23}(e_2 \wedge e_3)$$

For a 3D bivector,  $\vec{B} = B_{12}(e_1 \wedge e_2) + B_{13}(e_1 \wedge e_3) + B_{23}(e_2 \wedge e_3)$ , the components  $B_{ij}$  are the projections of  $\vec{B}$  onto the basis planes:

$$B_{12} = \vec{u}_1 \vec{v}_2 - \vec{u}_2 \vec{v}_1$$

$$B_{13} = \vec{u}_1 \vec{v}_3 - \vec{u}_3 \vec{v}_1$$

$$B_{23} = \vec{u}_2 \vec{v}_3 - \vec{u}_3 \vec{v}_2$$

Wait, you look familiar...

$$\begin{aligned} \vec{u} \wedge \vec{v} &= (\vec{u}_1 \vec{v}_2 - \vec{u}_2 \vec{v}_1)(e_1 \wedge e_2) \\ &\quad + (\vec{u}_1 \vec{v}_3 - \vec{u}_3 \vec{v}_1)(e_1 \wedge e_3) \\ &\quad + (\vec{u}_2 \vec{v}_3 - \vec{u}_3 \vec{v}_2)(e_2 \wedge e_3) \end{aligned}$$

$$\begin{aligned} \vec{u} \times \vec{v} &= (\vec{u}_1 \vec{v}_2 - \vec{u}_2 \vec{v}_1)e_3 \\ &\quad - (\vec{u}_1 \vec{v}_3 - \vec{u}_3 \vec{v}_1)e_2 \\ &\quad + (\vec{u}_2 \vec{v}_3 - \vec{u}_3 \vec{v}_2)e_1 \end{aligned}$$

Now introducing...

# The Geometric Product



# The Geometric Product

At the heart of geometric algebra is the geometric product:

$$\vec{u}\vec{v} = \vec{u} \cdot \vec{v} + \vec{u} \wedge \vec{v}$$

Similar to how an imaginary number has a real and imaginary part, the geometric product is a **multivector**; it has a scalar and bivector part.

Let  $\vec{u}$  be a vector, and let  $e_i, e_j$  be distinct basis vectors:

$$\vec{u}\vec{u} = \vec{u} \cdot \vec{u} + \vec{u} \wedge \vec{u} = \|\vec{u}\|^2$$

$$e_i e_j = e_i \cdot e_j + e_i \wedge e_j = e_i \wedge e_j = -e_j e_i$$

Let  $\vec{u}, \vec{v} \in \mathbb{R}^3$ . Explicitly, their geometric product is:

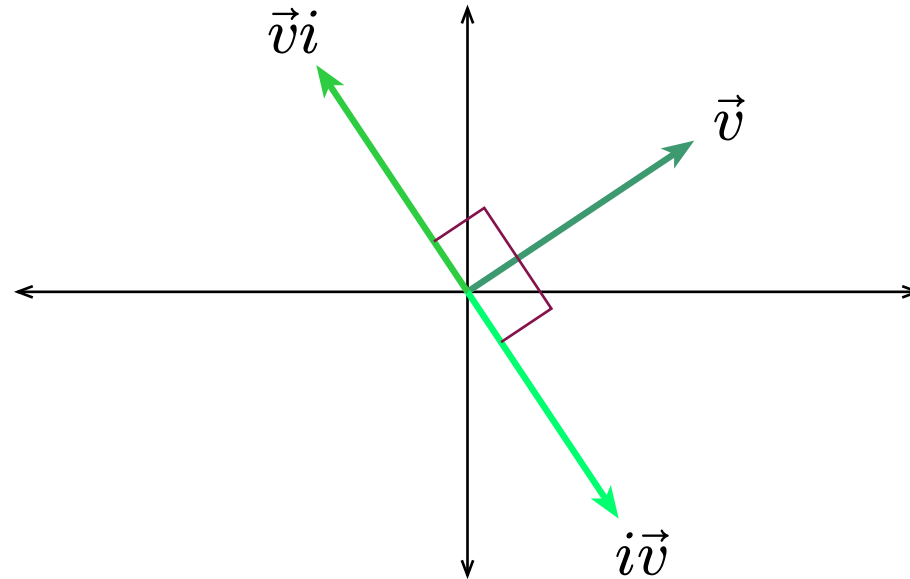
$$\begin{aligned} \vec{u}\vec{v} = & \vec{u}_1 \vec{v}_1 + \vec{u}_2 \vec{v}_2 + \vec{u}_3 \vec{v}_3 + (\vec{u}_1 \vec{v}_2 - \vec{u}_2 \vec{v}_1)(e_1 \wedge e_2) \\ & + (\vec{u}_1 \vec{v}_3 - \vec{u}_3 \vec{v}_1)(e_1 \wedge e_3) + (\vec{u}_2 \vec{v}_3 - \vec{u}_3 \vec{v}_2)(e_2 \wedge e_3) \end{aligned}$$

## Imaginary Numbers?

We return to  $\mathbb{R}^2$ . Let  $e_1, e_2$  be the basis vectors, and  $i := e_1 e_2$ .

Let  $\vec{v} = ae_1 + be_2$ . Then,

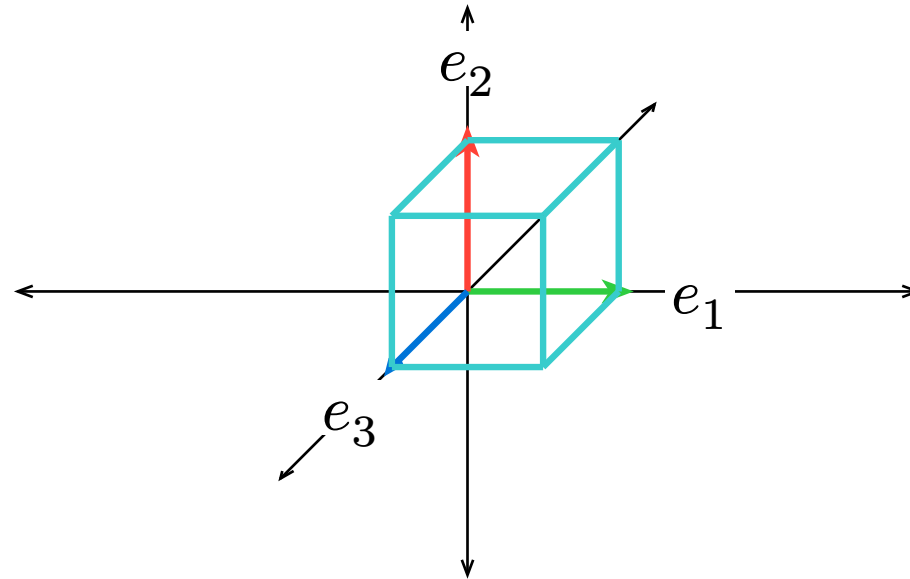
$$\vec{v}i = ae_2 - be_1, i\vec{v} = -ae_2 + be_1$$



Note that  $i^2 = (e_1 e_2)^2 = -e_2 e_1 e_1 e_2 = -1$ , just like imaginary numbers!

## Imaginary Numbers?

In  $\mathbb{R}^3$ , we have the orthonormal basis  $e_1, e_2, e_3$ , and  $i := e_1 e_2 e_3$ . This value is a trivector, and it looks the unit cube (with orientation).



Here,  $i$  commutes with any multivector  $\vec{A}$ . Note that  $e_1 i = e_2 e_3$  and  $e_1 = -e_2 e_3 i$ .

Remarkably,  $\vec{u} \wedge \vec{v} = i(\vec{u} \times \vec{v})$ .

## Quaternions?

If  $p, q$  are unit quaternions, then rotation of  $\vec{v}$  by  $pq$  is  $pq\vec{v}(pq)^{-1}$ . This looks a lot like rotation using rotors, but rotors are easier to understand.

Let  $x = e_1e_2, y = e_1e_3, z = e_2e_3$ :

	$x$	$y$	$z$
$x$	$-1$	$z$	$-y$
$y$	$-z$	$-1$	$x$
$z$	$y$	$-x$	$-1$

	$i$	$j$	$k$
$i$	$-1$	$k$	$-j$
$j$	$-k$	$-1$	$i$
$k$	$j$	$-i$	$-1$

$$(e_1e_2)^2 = (e_1e_3)^2 = (e_2e_3)^2 = (e_1e_2)(e_1e_3)(e_2e_3) = -1$$

$$i^2 = j^2 = k^2 = ijk = -1$$

## Quaternions?

In computer graphics, the geometric product can be used to encode 3D rotations. Let  $\vec{v}, \vec{a}, \vec{b}$  be linearly independent vectors. The reflection of  $\vec{v}$  about  $\vec{a}$  is given by:

$$\begin{aligned} R_{\vec{a}}(\vec{v}) &= \vec{v} - 2(\vec{v} \cdot \vec{a})\vec{a} \\ &= \vec{v} - 2\left(\frac{1}{2}(\vec{v}\vec{a} + \vec{a}\vec{v})\right)\vec{a} \\ &= \vec{v} - \vec{v}\vec{a}^2 - \vec{a}\vec{v}\vec{a} \\ &= -\vec{a}\vec{v}\vec{a} \end{aligned}$$

Likewise, reflecting about  $\vec{a}$  then about  $\vec{b}$  is given by:

$$R_{\vec{b}}(R_{\vec{a}}(\vec{v})) = -\vec{b}(-\vec{a}\vec{v}\vec{a})\vec{b} = \vec{b}\vec{a}\vec{v}\vec{a}\vec{b}$$

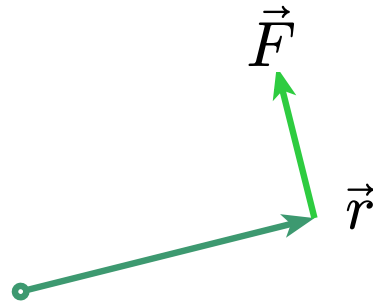
This encodes the rotation in the plane formed by  $\vec{a}, \vec{b}$  by  $2\angle\vec{a}\vec{b}$ . The product  $\vec{a}\vec{b}$  is called a **rotor**.

# Physical Applications

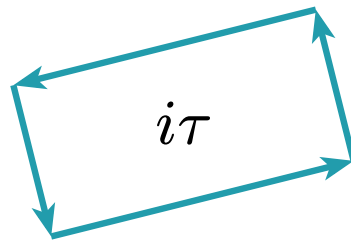
## Physics Made Intuitive

In 2D, a bivector is a **pseudoscalar**. Likewise, trivectors are 3D pseudoscalars. This is because the basis of bivectors in 2D has one element, and the basis of trivectors in 3D have one element.

For example, consider torque,  $\tau = \vec{r} \times \vec{F}$ :



For some reason,  $\tau$  points out of the page. But consider  $i\tau$ :



Which makes a lot more sense, in my opinion.

## Maxwell's Equations

Recall Maxwell's equations, which describe electromagnetism.

	$\vec{E}$	$\vec{B}$
$\cdot$	$\vec{\nabla} \cdot \vec{E} = \frac{\rho}{\epsilon_0}$	$\vec{\nabla} \cdot \vec{B} = 0$
$\times$	$\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$	$\vec{\nabla} \times \vec{B} = \mu_0 \left( \vec{J} + \epsilon_0 \frac{\partial \vec{E}}{\partial t} \right)$

$$\vec{\nabla} = \frac{\partial}{\partial x} \hat{x} + \frac{\partial}{\partial y} \hat{y} + \frac{\partial}{\partial z} \hat{z}$$

However, with geometric algebra, this can be made a lot simpler.



## Maxwell's Equations

We can add scalars and vectors now, so let's define a new gradient and some new variables.

$$\begin{aligned}\nabla &= \frac{1}{c} \frac{\partial}{\partial t} + \vec{\nabla} \\ &= \frac{1}{c} \frac{\partial}{\partial t} + \frac{\partial}{\partial x} \hat{x} + \frac{\partial}{\partial y} \hat{y} + \frac{\partial}{\partial z} \hat{z}\end{aligned}$$

$$J = c\rho - \vec{J}$$

$$F = \vec{E} + i\vec{B}$$

This gives us a new statement of Maxwell's equations...

$$\nabla F = J$$

# Generalizing to Higher Dimensions

## What happens when we go above 3D?

What is the geometric product of two bivectors?

$$(\text{Bivector})(\text{Bivector}) = \text{Scalar} + \text{Bivector} + 4\text{-vector}$$

This breaks our  $\vec{u}\vec{v} = \vec{u} \cdot \vec{v} + \vec{u} \wedge \vec{v}$  rule! Turns out, most of the time  $\vec{u}\vec{v} \neq \vec{u} \cdot \vec{v} + \vec{u} \wedge \vec{v}$ . In general:

$$\begin{aligned} (r\text{-vector})(s\text{-vector}) = & (|r - s|)\text{-vector} \\ & + (|r - s| + 2)\text{-vector} \\ & + (|r - s| + 4)\text{-vector} \\ & \vdots \\ & + (|r + s|)\text{-vector} \end{aligned}$$

This gives us the general form for the geometric product, between any vectors in any dimension.

Let's reflect...

## Let's reflect...

- In geometric algebra, you think about *oriented magnitudes*
- A lot of physical phenomena are pseudoscalars
- I lied to you guys in the abstract

Let  $V$  be a vector space with a symmetric bilinear form  $B : V \times V \rightarrow \mathbb{K}$ , where  $\mathbb{K}$  is a field. Say  $\otimes$  is the tensor product of  $u, v \in V$ , and consider the ideal  $\mathcal{J}(V, B)$  generated by elements of the form  $v \otimes w + w \otimes v - 2B(v, w)1$ . Then, consider  $T(V) := \bigoplus_{k \in \mathbb{Z}} T^k(V)$ , where  $T^k(V) = V \otimes V \dots \otimes V$  is the  $k$ -fold tensor product. Naturally, the Clifford (Geometric) algebra of  $V$  is the quotient:

$$\text{Cl}(V, B) = T(V) / \mathcal{J}(V, B)$$

Note that this is similar to the exterior algebra, which is  $T(V) / \mathcal{J}(V, B)$ , where  $\mathcal{J}(V, B)$  is the ideal generated by elements of the form  $v \otimes w + w \otimes v$ .

Thanks to the...



Mathematics Endowment Fund

**for funding my broke a\*\*.**